

**ANALYSIS OF THE EFFECT OF OXYGEN DOPING AND PRESSURE ON SUPERCONDUCTING TRANSITION TEMPERATURE IN METAL OXIDES***A. A. Kosov\***Mari State University  
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The two-orbital Hubbard model is used to obtain formulas for the fermion excitation spectrum in the energy bands hybridized by the Anderson interaction. An analysis of the lower part of the energy spectrum leads to a formula for the superconducting transition temperature  $T_c$  associated with the pairing of quasiparticles in one of the correlated bands. The dependence of  $T_c$  on pressure is analyzed, and the individual influence of carrier density enhancement and interaction strength is obtained as a function of oxygen concentration. The experimental discrimination made by Honma et al. [6] in  $Y_{0.9}Ca_{0.1}Ba_2Cu_3O_{7-\delta}$  by separating out the contributions due to carrier density and pairing strength can be reproduced quantitatively, and perhaps with further refinement, so can the carrier concentration. Although the prediction of the absolute value of the transition temperature using the present model is not accurate, it is clear that it furnishes a reasonably accurate description of the change in transition temperature with pressure. The component contributions due to the change in carrier concentration and due to the change in interaction strength as a function of oxygen concentration are also in reasonable agreement with the experimental results.

**1. INTRODUCTION**

The investigation of the mechanism of superconductivity in cuprate superconductors is related to the effects of oxygen doping and pressure on the superconducting transition temperature ( $T_c$ ) of these substances. One of the advantages of using high-pressure techniques is the ability to change atomic distances without substitution of components, which often causes some side effects [1, 2]. At the present time several review articles concerning high-pressure work in cuprate superconductors have been published. According to the experimental results and the conclusions of Shafer et al. [3] and Kubo et al. [4], the carrier concentration increases with increasing pressure in many high- $T_c$  materials. This increase in carriers is considered to be due to charge transfer from a charge reservoir layer to the Cu-O plane. The evidence for the dependence of the carrier concentration on pressure comes from measurements of the Hall number  $1/eR_H$  and the thermoelectric power under high pressure. Generally,  $T_c$  initially increases with increasing  $1/eR_H$ , but decreases when  $1/eR_H$  exceeds a specific value. However,

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the relation between  $T_c$  and  $1/eR_H$  in high-pressure experiments varies among different cuprate superconductors. For instance, in La-Sr-Cu-O ceramics,  $T_c$  increases with increasing pressure, but no change in  $1/eR_H$  is observed. On the other hand, in the case of Y-Ba-Cu-O, a variation of  $T_c$  with  $1/eR_H$  has been observed by several authors in high-pressure experiments. It was proposed [5] that the change in  $T_c$  due to pressure  $\Delta T_c$  should be expressed as the sum of two terms;  $(\Delta T_c)_c$  and  $(\Delta T_c)_p$ , where  $(\Delta T_c)_c$  is the change in  $T_c$  due to pressure-induced changes in the carrier density, and  $(\Delta T_c)_p$  is the change in  $T_c$  due to pressure-enhanced electron pairing (e.g., change in the electron-phonon coupling strength, or in the exchange coupling constant). Honma et al. [6] investigated the dependence of  $T_c$  on the Hall number by changing the oxygen content and the pressure in  $Y_{0.9}Ca_{0.1}Ba_2Cu_3O_y$ . They determined that the contribution of  $(\Delta T_c)_c$  to  $\Delta T_c$  increases with decreasing oxygen content.

In this work we use the idea that  $\Delta T_c = (\Delta T_c)_c + (\Delta T_c)_p$  and apply the Anderson-Hubbard two orbital model [7] to describe the experimental results of Ref. [6]. In Section 2 we introduce the Hamiltonian of the problem, the Green's functions of the quasiparticles in correlated bands, and the equation for  $T_c$ . Section 3 is devoted to the calculation of the pressure effect on the superconducting transition temperature. By proposing a simple relation between the value of pressure  $P$  and the width of the correlated band  $W$ , we can obtain the dependence of  $T_c$  on pressure. In addition, a comparison with experimental results on the dependence of  $T_c$ ,  $\Delta T_c$ ,  $(\Delta T_c)_c$  and  $(\Delta T_c)_p$  on the concentration of carriers in YCaBaCuO is made. Good agreement between the theoretical calculation of the dependence of  $(\Delta T_c)_c$  and  $(\Delta T_c)_p$  on pressure and the experimental results is found. It is concluded that the model under consideration is quite promising for studying the effects of oxygen doping and pressure on the superconducting transition in cuprate superconductors.

## 2. FORMULATION

One of the popular models used for describing a strongly correlated system is the Hubbard model [8]. Recently Kosov and Shilov [7, 9] studied the superconducting transition and pressure effects by using a unified Hamiltonian containing operators of the Hubbard two-orbital model and the Anderson interaction. The interaction considerably enhances the applicability of the Hubbard model and allowed the authors to describe the interaction of non-localized and localized electrons by proceeding from the mixing of their one-particle states.

The model in Refs. [7, 8] is based on the following Hamiltonian:

$$H = H_0 + H_{int} = \sum_i H_{0i} + \sum_{ijs} t_{ij} c_{is}^+ c_{is}, \quad (1)$$

$$H_{0i} = -\mu (n_{ia\uparrow} + n_{ia\downarrow} + n_{ic\uparrow} + n_{ic\downarrow}) + E (n_{ia\uparrow} + n_{ia\downarrow}) - \mathcal{H} (n_{ia\uparrow} - n_{ia\downarrow} + n_{ic\uparrow} - n_{ic\downarrow}) + \\ + I n_{ia\uparrow} n_{ia\downarrow} U (n_{ia\uparrow} + n_{ia\downarrow}) (n_{ic\uparrow} + n_{ic\downarrow}) + U_1 n_{ic\uparrow} n_{ic\downarrow} + V_0 (a_{i\uparrow}^+ c_{i\uparrow} + a_{i\downarrow}^+ c_{i\downarrow} + \text{H.c.}),$$

where  $c_{is}^+$ ,  $c_{is}$  and  $a_{is}^+$ ,  $a_{is}$  are field operators corresponding to free and localized electrons at the site  $i$  with spin projection  $s$ ;  $n_{ias} = a_{is}^+ a_{is}$  and  $n_{ics} = c_{is}^+ c_{is}$  are the operators for the number of electrons;  $\mu$  is the chemical potential;  $\mathcal{H}$  is the applied magnetic field,  $E$  is the one-particle energy of the  $a$ -electrons;  $I$ ,  $U$ , and  $U_1$  are the energy parameters defining intra-atomic correlation;  $I$  is the Hubbard interaction between localized electrons;  $U$  is the interorbital Coulomb interaction of  $c$ - and  $a$ -electrons;  $U_1$  is the repulsive interaction of  $c$ -electrons on

one site;  $V_0$  is the matrix element responsible for the hybridization of the  $c$ - and  $a$ -electronic states (Anderson's constant); and  $H_{int}$  describes the interstitial tunneling of  $c$ -electrons with transport integral  $t_{ij}$ .

We need to check our results according to separate condition:  $E > 0$  and  $E < 0$ . For this reason, we use the intermediate symbols:

$$Q = \begin{cases} B, & E > 0, \\ A, & E < 0, \end{cases} \quad S = \begin{cases} D, & E > 0, \\ C, & E < 0. \end{cases}$$

This means, for example, that for value  $R_K^S$  we have

$$R_K^S = \begin{cases} R_K^D, & E > 0, \\ R_K^C, & E < 0. \end{cases}$$

We assume that the density of states in the dispersion region has a rectangular shape [7-11]:

$$\rho(\varepsilon) = (1/2W)\theta(W^2 - \varepsilon^2),$$

where  $2W$  is the width of the  $c$ -band.

Carrying out calculations similar to those in Refs. [7, 11], we obtain the following expression for the chemical potential:

$$\mu = -\Delta/2 - WB_+/2 + 2(n_K + n_S)W(R_K^S)^2 - P_-/2, \quad (2)$$

$$P_{\pm} = (B_{\pm}^2 W^2 \pm 2\Delta_1 B_{\pm} W + \Delta_1^2)^{1/2},$$

$$B_{\pm} = \cos^2 \alpha (n_0 + n_Q) \pm (R_D^K)^2 (n_K + n_S).$$

Here  $n_0$ ,  $n_Q$ ,  $n_K$ , and  $n_S$  are the population densities of the energy states under investigation, which satisfy the condition  $n_0 + n_Q + n_K + n_S = 1$ ;

$$R_C^K = R_A^K = \cos \alpha (C_1 A_3 - A_1 C_3) / \sqrt{2} + \sin \alpha (C_1 A_2 - C_2 A_1),$$

$$A = (1 + A_{21}^2 + A_{31}^2)^{-1/2}, \quad A_2 = A_{21} A_1, \quad A_3 = A_{31} A_1,$$

$$A_{21} = (E_L - 2E - I)^2 / 2V_0^2, \quad A_{31} = (E_L - 2E - I)^2 / (E_L - U_1)^2,$$

$$B_2 = (1 + B_{12}^2 + B_{32}^2)^{-1/2}, \quad B_1 = B_{12} B_2, \quad B_3 = B_{32} B_2,$$

$$B_{12} = \frac{(E_M + 2E - U)(E_M - U_1) - 2V_0^2}{\sqrt{2} V_0 (E_M - U_1)}, \quad B_{32} = \frac{\sqrt{2} V_0}{E_M - U_1},$$

$$C_3 = (1 + C_{12}^2 + C_{23}^2)^{-1/2}, \quad C_1 = C_{13} C_3, \quad C_2 = C_{23} C_3,$$

$$C_{13} = \frac{(E_K - E - U)(E_K - U_1) - 2V_0^2}{2V_0^2}, \quad C_{23} = \frac{E_K - U_1}{\sqrt{2}V_0},$$

$$E_K < E_M < E_L.$$

The energy of the two-particle states  $E_K$ ,  $E_L$ , and  $E_M$  can be obtained by using the cubic equation

$$X^3 + AX^2 + BX + C = 0, \quad (3)$$

$$A = -(I + U + U_1 + 3E), \quad B = (I + 2E)(I + U) + U_1(3E + U + I) - 4(V_0)^2,$$

$$C = 2(V_0)^2(I + 2E + U_1) - U_1(E + U)(2E + I).$$

The roots  $X = \{X_m\}$  of Eq. (3) define the energy  $E_{K,L,M}$ :  $X_m = 2\mu + E_m$ ,  $m = K, L, M$ .

The chemical potential  $\mu$  is determined by the concentration of electrons in the dispersed correlated band, which equals  $n_c = (R_K^S)^2(n_K + n_S)$ . The dependence of  $n_c$  and  $\mu$  on  $V_0$  shows different behaviors in the cases  $E > 0$  and  $E < 0$ . An increase in the hybridization parameter  $V_0$  leads to a decrease in  $n_c$  for  $E > 0$ , and to an increase for  $E < 0$ .

A transition to the Hubbard operators allows the use of the Green's temperature function technique to take the interstitial jump term into account in order to study the superconducting properties of the model. An analysis of the lower part of the energy spectrum leads to the following formula for the superconducting transition temperature associated with the pairing of quasiparticles in one of the correlated bands:

$$\frac{T_c}{2W} = 0.57 \left[ -\frac{\xi_{10}(-W)\xi_{10}(W)}{W^2} \right]^{1/2} \exp \left[ -\frac{1}{A(n, t_0)} \right], \quad (4)$$

$$\xi_{10}(-W) = -2(R_K^S)^2(n_K + n_S),$$

$$\xi_{10}(W) = -\mu + (B_+W - P_+ - \Delta)/2,$$

$$A(n, t_0) = \Gamma(n, t_0)/\Lambda(n, t_0),$$

$$\Gamma(n, t_0) = t_0 [\cos^2 \alpha(n_0 + n_Q)(E_S - E_K)(E_S + B_-t_0) + (R_K^S)^2(n_K + n_S)E_Q(E_Q - E_K + B_-t_0)],$$

$$\Lambda(n, t_0) = W [B_+(2\mu + \Delta) + B_-\Delta_1] (2\mu + \Delta - B_+t_0).$$

The quantity  $A(n, t_0)$  plays the role of the quasiparticle scattering amplitude with different spin orientation. The attraction between quasiparticles in a correlated band takes place under the conditions:

$$A(n, t_0) > 0, \quad -W < t_0 < W, \quad \xi_{10}(W) \geq 0. \quad (5)$$

The conditions (5) can be used to determine the concentration  $n_{S1} < n < n_{S2}$  for which  $T_c \neq 0$  holds. Solving the equation  $\xi_{10}(W) = 0$  gives the following value for  $n_{S2}$ :

$$n_{S2} = \frac{2 \cos^2 \alpha}{\cos^2 \alpha + (R_K^S)^2}. \quad (5a)$$

The condition  $A(n, t_0) = 0$  gives the following result for  $n_{S1}$ :

$$n_{S1} = \frac{2 \cos^2 \alpha}{\cos^2 \alpha + 2(R_K^S)^2}. \quad (5b)$$

In particular, if we put  $V_0 = 0$  and  $E > 0$  in (5a) and (5b), we obtain  $n_{S1} = 2/3$  and  $n_{S2} = 1$ . This result has been obtained by Zaitsev and Ivanov [12] in the framework of the one-orbital Hubbard model (the so-called «kinematic mechanism of superconductivity»).

### 3. PRESSURE EFFECTS

In applying the two-orbital Anderson-Hubbard model to describe the pressure dependence of  $T_c$ , we choose to examine the energy parameters, which more sensitively depend on the value of pressure. Since  $U$ ,  $I$ ,  $U_i$ ,  $V_0$ , and  $E$  are on-site properties, their pressure dependence can be neglected. The transport integral  $t_p$  depends on the spatial distribution of atoms and is changed by applied pressure. Let us consider the region where  $W$  depends linearly on  $t_p$  ( $W \propto t_p$ ). According to Marsiglio and Hirsh [13], the transport integral in cuprate superconductors can be expressed through the lattice parameters of the  $\text{CuO}_2$  planes  $a$ ,  $b$ , and  $c$  by the formulas

$$t_{\parallel} = \frac{\hbar^2}{2m_{\parallel}a}, \quad t_{\perp} = \frac{\hbar^2}{2m_{\perp}c},$$

where  $m_{\parallel}$  and  $m_{\perp}$  are the respective effective masses.

Neglecting any pressure dependence of the effective masses, we estimate the magnitude of  $dW/dP$ :

$$\frac{dW}{dP} \propto \frac{dt_p}{dP} \sim -2W \frac{d \ln a}{dP} = 2W k_a. \quad (6)$$

Here  $k_a$ ,  $k_b$ , and  $k_c$  are the compressibility components along each crystallographic direction, defined by

$$k_a = -\frac{d \ln a}{dP}, \quad k_b = -\frac{d \ln b}{dP}, \quad k_c = -\frac{d \ln c}{dP}.$$

In order to simplify the numerical estimation of the results, we consider the case  $k_a = k_b = k_c$ . Using the relations in the equation (6), we obtain the following expression for the dependence of the bandwidth on pressure

$$W(P) = W(P=0) \exp(2k_a P). \quad (7)$$

Formulas (4) and (7) allow us to express  $T_c$  as a function of pressure. Using expression (6) we can obtain formulas for  $dT_c/dP$  and  $d \ln T_c/dP$ :

$$\frac{dT_c}{dP} = T_c \frac{d \ln T_c}{dP} \sim 2W k_a \frac{dT_c}{dW}. \quad (8)$$

$$\frac{dT_c}{dW} = \frac{T_c}{2W} \left\{ 1 + 2W \left[ B_- + \frac{B_+^2(P_+ - P_-)W - \Delta_1 B_- (P_+ + P_-)}{2P_+ P_-} \right] + \frac{2W}{A^2(n, t_0)} \frac{dA(n, t_0)}{dW} \right\},$$

$$\frac{dA(n, t_0)}{dW} = -\frac{A(n, t_0)}{W} + \frac{1}{\Lambda(n, t_0)} \frac{d\mu}{dW} \{ [2(2\mu + \Delta) - B_+ t_0] \times$$

$$\times [\cos^2 \alpha(n_0 + n_Q)(E_S - E_K)(E_S + 2B_- t_0) + (R_K^S)^2(n_K + n_S) \times$$

$$\times E_Q(E_Q - E_K + 2B_- t_0) + A(n, t_0)B_+ W (B_+(2\mu + \Delta) + \Delta_1 B_-)] -$$

$$- 2WA(n, t_0) [2B_+(2\mu + \Delta) + \Delta_1 B_- - B_+^2 t_0] \},$$

$$\frac{d\mu}{dW} = 2(R_K^S)^2(n_K + n_S) - \frac{B_+}{2} - \frac{WB_+^2 - \Delta_1 B_-}{2P_-}.$$

By means of formula (8) we obtain a theoretical expression for  $(\Delta T_c)_p$ :

$$(\Delta T_c)_p = \frac{dT_c}{dP} \Delta P = \frac{dT_c}{dW} \Delta W. \tag{9}$$

To compare our results with the experimental data of Ref. [6] we obtain the value  $(\Delta T_c)_c$  from formula (4):

$$(\Delta T_c)_c = \frac{dT_c}{dn} \Delta n = \frac{dT_c}{dn} \frac{dn}{dP} \Delta P = \frac{dT_c}{dn} \frac{dn}{dW} \Delta W. \tag{10}$$

#### 4. DISCUSSION

We have chosen to compare the theoretical results presented in the previous section with the recent experimental results of Honma et al. [6], who studied the pressure dependence and the effect of oxygen doping on the transition temperature of the  $Y_{0.9}Ca_{0.1}Ba_2Cu_3O_y$  (YCBCO) system. These authors systematically investigated the critical temperature dependence and the dependence of the Hall number  $1/eR_H$ , on pressure, in order to be able to distinguish between those effects which result from changes in carrier density  $(\Delta T_c)_c$ , and changes in the coupling strength (electron pairing)  $(\Delta T_c)_p$ .

The first comparison we shall examine is the dependence of the critical temperature  $T_c$ , on the oxygen content  $y$ . The theoretical results for a pressure of 2 GPa and a compressibility  $k_a$  of  $0.0024 \text{ GPa}^{-1}$  are displayed in Fig. 1, where five curves are plotted, each with a variation in one or more of the following bandwidth-normalized parameters: the single-particle energy  $E/W$ , the interatomic correlation energies  $I/W$ ,  $U_1/W$ , and  $U/W$ ; and the hybridization energy  $V_0/W$ . The values of the parameters used in the various curves used to fit the data are listed in the table below.

Table

Parameter values used

Curve	$-E/W$	$I/W$	$U_1/W$	$U/W$	$V_0/W$
1	-1.9	6.0	3.5	2.4	1.8
2	-1.9	6.0	3.5	2.4	2.0
3	-1.7	6.0	3.5	2.4	1.8
4	-1.8	5.0	2.5	2.8	1.8
5	-1.9	6.0	2.4	3.5	2.0

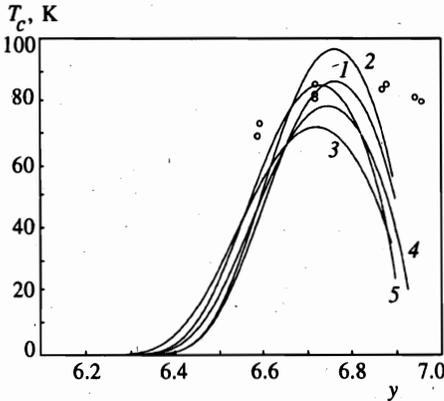


Fig. 1. Dependence of the function  $T_c$  on the oxygen content  $y$  for  $Y_{0.9}Ca_{0.1}Ba_2Cu_3O_y$  obtained for the parameters  $k_a = 0.00024 \text{ GPa}^{-1}$ ,  $P = 2 \text{ GPa}$

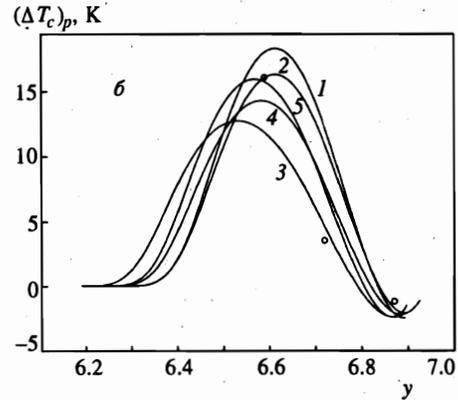
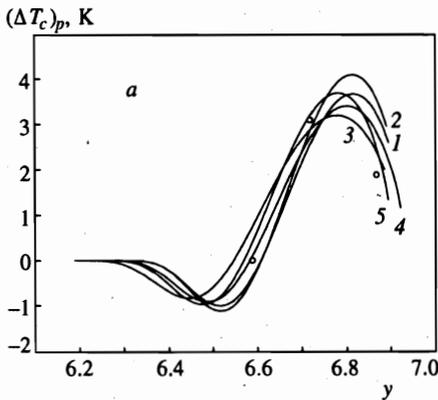


Fig. 2. Dependences of  $(\Delta T_c)_p$  (a) and  $(\Delta T_c)_c$  (b) on the oxygen content  $y$  for  $Y_{0.9}Ca_{0.1}Ba_2Cu_3O_y$  obtained for the parameters  $k_a = 0.00024 \text{ GPa}^{-1}$ ,  $P = 2 \text{ GPa}$

The data of Fig. 1 in the paper by Honma et al. are represented by open circles. As can be seen from the figure, the general concave-downward form of the experimental results can be qualitatively reproduced by the theoretical curves with the same maximum value, and with little variation exhibited for different parameter combinations. The requisite width of the curve exhibited by the experimental data, however, is not readily attained with any reasonable variation in the parameters. The difficulty in fitting the theory to an absolute quantity like the critical temperature is not unexpected, since the absolute values of the potentials chosen are involved in making such a comparison.

In Figs. 2a and 2b are plotted the critical temperature change at a pressure of 2 GPa and a compressibility  $k_a$  of  $0.0024 \text{ GPa}^{-1}$ , due to changes in carrier density and pairing strength  $(\Delta T_c)_c$  and  $(\Delta T_c)_p$ , respectively, as a function of oxygen content  $y$ . Here, the agreement with the experimental results of Honma et al (open circles) is fairly good. Once again, the theoretical curves exhibit a concave-downward behavior over the range of  $y$  that was measured. Of the parameter sets chosen for illustration, it appears that the parameters corresponding to curve 5 ( $-E/W = -1.9$ ,  $I/W = 6.0$ ,  $U_1/W = 2.4$ ,  $U/W = 3.5$ , and  $V_0/W = 2.0$ ) give the closest fit to the experimental data in both cases. Figure 3 illustrates the sum of the two effects to give the total change in temperature,  $(\Delta T_c)_c + (\Delta T_c)_p = \Delta T_c$  plotted vs. oxygen content  $y$ ,

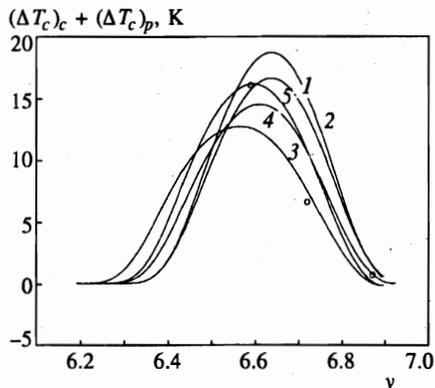


Fig. 3.

Fig. 3. Dependence of  $(\Delta T_c)_p + (\Delta T_c)_c$  on the oxygen content  $y$  for  $Y_{0.9}Ca_{0.1}Ba_2Cu_3O_y$  obtained for the parameters  $k_a = 0.00024 \text{ GPa}^{-1}$ ,  $P = 2 \text{ GPa}$

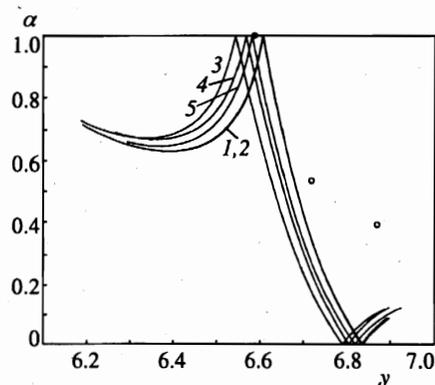


Fig. 4

Fig. 4. Dependence of the ratio  $\alpha = |(\Delta T_c)_c| / (|(\Delta T_c)_p| + |(\Delta T_c)_c|)$  on the oxygen content  $y$  for  $Y_{0.9}Ca_{0.1}Ba_2Cu_3O_y$  obtained for the parameters  $k_a = 0.00024 \text{ GPa}^{-1}$ ,  $P = 2 \text{ GPa}$

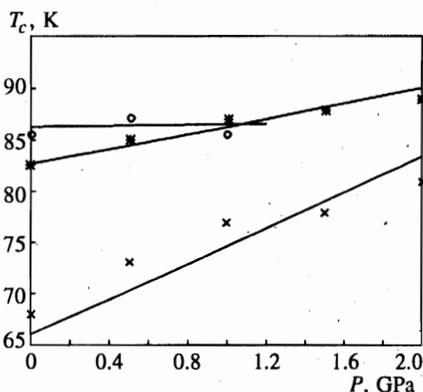


Fig. 5.  $T_c$  as a function of pressure in  $Y_{0.9}Ca_{0.1}Ba_2Cu_3O_y$  ( $\times$  —  $y = 6.59$ ,  $*$  —  $y = 6.72$ ,  $o$  —  $y = 6.87$  — experimental data [6]) obtained for the parameter  $k_a = 0.00024 \text{ GPa}^{-1}$  (used for curve 5)

along with the experimental data. Once again the theoretical treatment yields a set of concave-downward curves which qualitatively approximate the data with the best fit to the experimental data given by curve 5. Overall, the theoretical model appears to quite accurately account for the observed experimental variation with oxygen content. In contrast with the first plot, these comparisons involve the change in critical temperature with pressure and so can be more realistically accounted for by the theoretical model.

For the Hall number  $1/eR_H$ , which should scale with the ratio of the carrier density enhancement of the critical temperature  $(\Delta T_c)_c$  to the sum of the absolute values of the individual critical temperature enhancements  $|(\Delta T_c)_c| + |(\Delta T_c)_p|$ , the plot vs. oxygen concentration  $y$  shown in Fig. 4 illustrates the comparison between theory and experimental data of Honma et al., who labelled this ratio  $\alpha$ . Here the fit is not good with any of the parameter variations that were tried. Although the experimental data fall on curve 4 near

$\Delta T_c / (|(\Delta T_c)_c| + |(\Delta T_c)_p|) = 1$ , where the temperature change is almost entirely due to change in carrier concentration, the slopes are clearly not in agreement. We believe this discrepancy arises from the reduced accuracy in evaluating this fraction when the Hall coefficient changes sign, which it does in the case where  $(y, \alpha) = (6.87, 0.39)$ . The theoretical curve nevertheless gives a reasonable qualitative description of the variation in this parameter, as a decreasing function of oxygen content. In order to determine how  $dT_c/dn$  varies with oxygen concentration  $y$ , we have used a third-order polynomial fit of the form:  $dn(y) = a_3y^3 + a_2y^2 + a_1y + a_0$ . The fitting parameters are obtained by using the data of Honma et al. at  $y = 6.59, 6.72$  and  $6.87$ , respectively, and by setting  $dn(7.0) = 0$ . The values of the fitting parameters are  $a_0 = -461.1850$ ,  $a_1 = 141.5920$ ,  $a_2 = -11.6100$ ,  $a_3 = 0.1135$ . This relation provides a reasonable idea about the strength of the effect of charge carrier density upon  $\Delta T_c$ .

Finally, in Fig. 5, we plot the total change in the transition temperature vs. pressure as obtained from the theoretical model. The experimental results are shown as data points according to the legend. It is apparent that the quality of the agreement is good, with no more than 10% discrepancy. The rate of increase in  $T_c$  with pressure is faithfully reproduced for all three fractional oxygen contents studied by Honma et al. It is important to note the nearly zero slope at the highest concentration,  $y = 6.87$ .

## 5. CONCLUSIONS

In this paper we have demonstrated that the application of the Hubbard model provides a good basis for describing the observed variation of the pressure dependence of the critical temperature on oxygen concentration in the  $Y_{0.9}Ca_{0.1}Ba_2Cu_3O_y$  (YBCO) system. It appears that the experimental discrimination made by Honma et al. in separating out the contributions due to carrier density and pairing strength can be reproduced quantitatively, and perhaps with further refinement, as can the carrier concentration. Although the prediction of the absolute value of the transition temperature is not accurate using the present model, it is clear that it furnishes a reasonably accurate description of the change in transition temperature with pressure. The component contributions due to change in carrier concentration and due to change in interaction strength as a function of oxygen concentration are also in reasonable agreement with the experimental results.

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