SURFACE TENSION OF He³-He⁴ SOLUTIONS

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The surface tension α of liquid mixtures of helium isotopes, containing 9.5, 15.0, 19.0, 50.0 and 75.7% He³, was measured in the 1.3-4.2°K temperature range. The value of α decreased smoothly with increase of the He³ content, and the values were lower than those expected on the assumption of additivity. The results are compared with the conclusions of Prigogine's theory. ^[14]

HE temperature dependence of the surface tension of pure helium isotopes has been studied in adequate detail. [1-6] Measurements of the surface tension of He³-He⁴ solutions (up to 3% He³) have shown [7] that the He³ impurity lowers the surface tension of He⁴.

In the present work the surface tension measurements were extended to higher concentrations of He³ (up to 75%). We used a measuring device (Fig. 1b) and apparatus for observations (Fig. 1a) similar to those used earlier. ^[7] In the present case the volume of the device was about 0.2 cm³, the diameter of the wide tube was 2.89 mm and the capillaries were of diameters from 0.12 to 0.22 mm. To stir the mixture in order to reach equilibrium more quickly, we used an iron rod which could be moved vertically inside the wide tube and which followed the motion of a permanent magnet next to the device. Reciprocating motion was transmitted to the magnet by a wire joining it to an SD-60 motor with a cam on its axle.

In the region of the existence of He II we used a controller ^[8] which kept the temperature constant to within 5×10^{-6} deg K. The difference of the heights of the liquid levels in the "knee" of the device was measured with a cathetometer KM-5 to within 0.01 mm.

Before and after the measurements on heliumisotope solutions we carried out control experiments on pure He⁴. The results obtained in these experiments agreed well with the published data.^[1,2]

The surface tension was calculated from

$$2\alpha (1/b_1 - 1/b_2) = (\rho_l - \rho_v) gh,$$

where α is the surface tension; b_1 and b_2 are the radii of curvature of the lowest points of the menisci in two tubes of radii r_1 and r_2 ; ρ_l and ρ_v are, respectively, the densities of the liquid and the vapor; g is the acceleration due to gravity; h is the height between the lowest points of the menisci. The cor-





rection for the difference between the radii of curvature of the menisci (b_1 and b_2) and the radii of the capillary (r_1) and the tube (r_2) was calculated by the method of successive approximations using Sugden's tables, ^[9] and amounted to $\approx 5\%$. The density of the liquid was determined from the experimental data ^[10] and the density of the vapor from the known vapor densities of the pure components, ^[11,12] on the assumption that they are additive. The concentration of He³ in the vapor was calculated from the data in ^[13]. The accuracy of the measurements at all temperatures and concentrations was not less than 4% and for concentrations up to 20% it was higher (about 1%).



FIG. 2. Temperature dependence of the surface tension of He⁴, He³, and their solutions: $\Delta = \begin{bmatrix} 1 \\ 2 \end{bmatrix}$; $\bullet = \begin{bmatrix} 3,7 \\ 2 \end{bmatrix}$; $\diamond = \begin{bmatrix} 4,5 \\ 2 \end{bmatrix}$; $\Box = \begin{bmatrix} 6 \\ 2 \end{bmatrix}$; $\circ =$ results of present work. Curve a represents a solution containing 9.5% He³, b = 15.0%, c = 19.0%, d = 50.0%, e = 75.7%.

The temperature dependence of the surface tension of the pure helium isotopes and their solutions is shown in Fig. 2. The uppermost curve represents pure He⁴, the lowest one pure He³. The three curves between the uppermost one and curve a represent the temperature dependence of the surface tension of solutions containing 0.8, 1.9, and 3.0% He³ respectively.^[7] Figure 2 indicates that the surface tension decreases with increase of the concentration of the light isotope and with increase of temperature.

From the curves of the temperature dependence of the surface tension for various concentrations we plotted isotherms, two of which are shown in Fig. 3. These curves clearly show the departure from additivity (represented by the dashed line) in the direction of lower values of α .



FIG. 3. Dependence of the surface tension on the concentration of the light isotope: 1) T = 1.5° K; 2) T = 2.1° K.

Confining ourselves to the second terms of the expansion we find that the excess surface tension

$$\alpha_e = \alpha - \alpha_3 X - \alpha_4 (1 - X) \tag{1}$$

where α , α_3 and α_4 are the surface tensions of the solution and of He³ and He⁴, respectively, and X is the molar concentration of He³ in the solution, is given by the following expression:

$$\alpha_e = -\frac{X\left(1-X\right)}{2} \frac{\partial^2 \alpha}{\partial \left(\Lambda^2\right)^2} (\Lambda_3^2 - \Lambda_4^2)^2 - \frac{X\left(1-X\right)\omega}{2kT} (\alpha_4 - \alpha_3)^2, \qquad (2)$$

where $\Lambda = h/d \sqrt{m\epsilon}$ is a nondimensional parameter; $\Lambda_4 = 2.64$ (for He⁴); $\Lambda_3 = 3.05$ (for He³); ω is the surface per single particle.

Figure 4 shows the dependence of the excess surface tension on the concentration of He^3 . It is clear that the theory gives a qualitatively correct description of the phenomenon and that it agrees quantitatively with the experimental results at low He^3 concentrations.



FIG. 4. Dependence of the excess surface tension on the concentration of He³: the continuous curve represents the experimental results according to the definition of Eq.(1); the dashed curve is theoretical, calculated from Eq.(2).

It is our intention to extend the measurements of the surface tension toward lower temperatures.

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