Reactions	Au ¹⁹⁷ + O ¹⁶		Au ¹⁹⁷ + N ¹⁴	Au ¹⁹⁷ + C ¹²	
Ion energies, Mev l^2_{max}	1 02 2700	85 760	89 2060	78 1750	64 530
Emission ratio of Cd ¹¹⁵ and Cd ^{115^m}	0.43	0.50	0,47	0,55	0.64

method after careful purification. The purity of the solution of cadmium was monitored by the β -particle half-life and energy. The measurement of the radiation intensity was taken by an end-window Geiger counter (MSG-17). The thickness of the CdS layer deposited from the filter paper onto a backing was $\sim 10 \text{ mg/cm}^2$. The variation of the thickness of the layer of CdS from experiment to experiment, $2-3 \text{ mg/cm}^2$, need not have a considerable effect on the scattering and self-absorption of β radiation.³ It is also necessary to note that in all our experiments, apparently, a considerable quantity of Ag^{115} ($T_{1/2} = 20$ min.) was formed, which decays $90\%^4$ to Cd^{115} by β emission. After irradiation, the foil was kept intact for a sufficiently long time for the full decay of Ag¹¹⁵. Additional experiments showed that approximately 30% of Cd¹¹⁵ was formed as a result of Ag¹¹⁵ decay.

The results of the principal experiments are presented in the table (the given data are uncorrected for the formation of Cd^{115} from Ag^{115}).

The table lists also the computed l_{\max}^2 for the

LATERAL DISTRIBUTION FUNCTION OF THE FLUX OF CHARGED PARTICLES IN AN INDIVIDUAL EXTENSIVE ATMOS-PHERIC SHOWER

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 - Submitted to JETP editor October 14, 1959
 - J. Exptl. Theoret. Phys. (U.S.S.R.) 38, 297-298 (January, 1960)

L HE apparatus in operation at the Moscow State University for an all-out investigation of extensive atmospheric showers (EAS) makes it possible to investigate the individual characteristics of each recorded shower. In this note we report experimental data on the lateral distribution functions corresponding reactions (l_{max}^2) is the maximum orbital momentum of the ion at which fusion of the impinging nuclei arises). It is apparent that with a change of l_{max}^2 from 500 to 3000, the emission ratio of Cd¹¹⁵ and Cd^{115m} changes very slightly.

In conclusion we consider it our pleasant duty to thank corresponding member of the Academy of Sciences U.S.S.R. G. N. Flerov for guidance in this present work.

¹ V. M. Strutinskiĭ, Атомная энергия (Atomic Energy) **6**, 508 (1957).

² P. Kruger and N. Sugarman, Phys. Rev. **99**, 1468 (1955).

³H. G. Hicks and R. S. Gilbert, Phys. Rev. **120**, 1286 (1958).

⁴ Hollander, Perlman, and Seaborg, Revs. Modern Phys. **25**, 534 (1953).

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of the flux density of charged particles up to 25 meters from the shower axis. To obtain these data, we used a hodoscope with a large number of Geiger-Müller counters and a core detector consisting of ionization chambers arranged in two rows.¹ The first row of ionization chambers, shielded with lead, is used to determine the distribution of the energy flux carried by the electronphoton component near the shower axis. A direct examination of the distribution of the energy flux in the first row of ionization chambers makes it possible to determine the position of the shower axis for showers with a sufficient number of particles, with accuracy on the order of the chamber dimensions (25 cm). The hodoscopic counters were used to determine the flux density of the charged particles at different distances from the shower axis. To investigate the lateral-distribution functions of the flux density of the charged particles in an individual shower, we selected 26 of the densest showers (with $N \ge 10^5$), whose axes fall in the first row of the ionization chambers, for only in such showers can the charged-

N	S								
NO.	0.8	0.9	1.0	1.1	1.2	1.3	1.4	1.5	
$\begin{array}{c} 1\\ 2\\ 3\\ 4\\ 5\\ 6\\ 7\\ 8\\ 9\\ 10\\ 11\\ 12\\ 13\\ 14\\ 15\\ 16\\ 17\\ 18\\ 19\\ 20\\ 21\\ 22\\ 24\\ 25\\ 26\end{array}$	0.4 ≪0.01 0.01	$\begin{array}{c} 0.1 \\ 0.05 \\ 0.7 \\ 0.2 \\ 0.01 \\ 0.05 \end{array}$	$ \begin{array}{c} \ll 0.01 \\ 0.4 \\ 0.6 \\ 0.1 \\ 0.5 \\ 0.01 \\ 0.01 \\ \ll 0.01 \end{array} $	$\begin{array}{c} 0.15\\ 0.01\\ 0.05\\ 0.15\\ 0.5\\ 0.3\\ 0.5\\ 0.01\\ 0.35\\ < 0.01\\ 0.03\\ 0.05\\ \end{array}$	$\begin{array}{c} 0.01\\ 0.02\\ 0.1\\ 0.25\\ 0.75\\ 0.15\\ 0.25\\ < 0.01\\ 0.15\\ 0.15\\ 0.15\\ 0.15\\ 0.15\\ 0.01\\ 0.15\\ 0.1\\ \ll 0.01\\ \end{array}$	$\begin{array}{c} 0.2\\ 0.05\\ 0.15\\ < 0.01\\ < 0.01\\ < 0.01\\ 0.4\\ 0.6\\ 0.9\\ 0.45\\ 0.1\\ 0.4\\ 0.5\\ 0.2\\ 0.3\\ < 0.01\\ \end{array}$	$< \stackrel{0.01}{\stackrel{0.01}{_{0.01}}} \\ < \stackrel{0.01}{_{0.01}} \\ < \stackrel{0.01}{_{0.02}} \\ \stackrel{0.02}{_{0.05}} \\ \stackrel{0.1}{_{0.8}} \\ \stackrel{0.4}{_{0.4}} $	$\begin{array}{c} 0.1 \\ 0.02 \end{array}$	

particle flux density be investigated in this range of distances with good accuracy.

The experimental charged-particle lateral distribution functions obtained for each of the selected showers were compared with the theoretical functions, calculated by Nishimura and Kamata² for various values of the cascade parameter S. The theoretical curves were normalized here to the number of particles experimentally observed in a circle of radius 25 m. The Pierson matching criterion was used to choose the theoretical curve corresponding to the experimental data. The results [the Pierson function $P(\chi^2)$] are listed in the table, which shows which values of the parameter S characterize the charged-particle flux density lateral distribution functions in the registered showers with different particle numbers N.

The experimental data given indicate the existence of extensive atmospheric showers of various ages near sea level.

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A MEASUREMENT OF THE SURFACE TEN-SION AT THE BOUNDARY BETWEEN THE SUPERCONDUCTING AND NORMAL PHASES OF INDIUM

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The value of the surface tension σ_{ns} at the boundary between the superconducting and normal phases of indium has been measured by a method described previously.¹

A single-crystal disc with diameter 50 mm and thickness 2.06 mm, fabricated of indium with impurity content ~0.002%, was placed in a magnetic field directed at an angle of 15° to the specimen surface. The structures observed were completely analogous to those observed in tin¹ at the same values of H/H_c . The period of the structure was measured for various fields and temperatures and the quantity $\Delta(T) = \sigma_{\rm ns} (8\pi/H_c^2)$ was calculated.*

In the figure are given the results of measurements on Δ made in the range 2.11 to 3.245°K with an accuracy of 8 - 10%. In this range the results can be described within the limits of error by the relationship

¹Vernov, Goryunov, Zatsepin, Kulikov, Nechin, Strugal'skiĭ, and Khristiansen, JETP **36**, 669 (1959), Soviet Phys. JETP **9**, 468 (1959).

² J. G. Wilson, ed. <u>Progress in Cosmic Ray</u> <u>Physics</u>, (Russ. Transl.), vol. 3, 1959, p. 7.