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Translated by D. C. West 307

Measurement of the Absolute Quantum Yield of Photoluminescence of Alkali-Halide Crystals

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I N earlier works the dependence of the relative quantum yield of photoluminescence of crystals of KI-Tl and NaI-Tl on the concentration of the activator¹ and the wavelengths of the exciting light² was investigated. In the present work the absolute quantum yield of luminescence of these crystals during excitation in the absorption region of the activator was investigated.

For determining the absolute quantum yield a method was used in which the number of luminescent quanta and the number of absorbed quanta of exciting light was measured with the aid of a substance which is luminescent under the action of the one as well as the other radiation, and which has a constant quantum yield in the entire measurable region. This method has been applied earlier to the measurement of the absolute luminescence yield of powdered phosphors which radiate in the ultra-violet³, and also of some finely-crystalline organic phosphors⁴.

Our measurements were carried out in a photometric sphere, which made it possible to take into account the entire flux of luminescence. As a measurable luminescent substance a powdered, light yellow phosphor was taken. The preliminary tests showed that the quantum luminescence yield of this phosphor does not depend on the wavelength in the region $240-500 \,\mathrm{m}\mu$, and the -bsorption coefficient of the surface of the powdered layer of phosphor is constant (about 0.95) in this region. The excitation and radiation spectrum of the crystals KCl-Tl, KI-Tl, and NaI-Tl is in this region.

The measurements were carried out with light excitation from a hydrogen lamp through a spectrometer SF-4, with a cap enclosing an FEY-19 photomultiplier and an amplifier*. The crystal was placed in an aluminum sphere (ϕ 4cm), which had a slot $(4 \times 10 \text{ mm})$ for the entrance of the exciting light, and an aperture $(\phi 10 \text{ mm})$ for the exit of the luminescent light. The sphere was lined in the interior with a thick layer of phosphor. The exciting light fell on the crystal and part of it was reflected while the remainder was absorbed, producing luminescence of the crystal (special attention was given so that the absorption of exciting light in the crystal was complete). The reflected light as well as the luminescent light of the crystal caused a glow of the phosphor, which was registered by the photomultiplier.

The reflected (Fresnel) exciting light was measured with the help of another nonabsorbing crystal, the back face of which was blackened. A crystal of NaCl was used for this purpose. The small difference in the index of refraction is unimportant, in view of the small size of the correction. The incident exciting light (in order to preserve the geometry) was similarly measured with the aid of a nonabsorbing crystal, the back face of which in this case was covered with magnesium oxide. It is obvious that if in the first measurement we obtain the reading a, in the second b, and in the third c, then the quantum yield will equal $\eta = (a - b) / (c - d)$. The results of the measurements of yield are presented in the following Table.

1		
9/6	0.80	
240	0.80	
245	0.70	Different crystals One and the same crystal, different method of calculation of the reflected light
285	$\left. \begin{array}{c} 0.77 \\ 0.74 \\ 0.76 \end{array} \right\}$	
260 2 9 3	0.40 0.61	
	246 245 285 260 293	$\begin{array}{cccc} 246 & 0.80 \\ 0.80 \\ 0.70 \\ 0.93 \end{array} \\ \begin{array}{c} 285 \\ 0.77 \\ 0.74 \\ 0.76 \\ 0.76 \\ 260 \\ 293 \end{array} \\ \begin{array}{c} 0.77 \\ 0.76 \\ 0.40 \\ 293 \end{array} \\ \begin{array}{c} 0.61 \end{array} \end{array}$

The concentration of Tl in the crystals KI-Tl and NaI-Tl was sufficient so that during the excitation in the shortwave region of absorption the Tl yield no longer depended on its concentration¹. In KCl-Tl, according to our measurements, the concentration of Tl was 1.5×10^{-4} gm Tl/gmKCl; the data for KCl-Tl correspond well with those available in the literature⁵. It should be noted that in KI-Tl the variations in the value of the yield for various crystals (with high content of Tl) is found to be greater during excitation with $\lambda = 245 \text{ m}\mu$ than with $\lambda = 285 \text{ m}\mu$. This is easily explained by the different degree of overlapping of the short wave region of Tl absorption with the boundary of absorption of the lattice itself, which, as has been shown², is inactive.

The above method was also used to measure the quantum yield of several organic monocrystals. The coefficient of reflection in this case was measured directly. With excitation of $\lambda = 250 \text{ m}\mu$ the yield amounts to: for tolane 0.57; for dibenzyl 0.44; for stilbene 0.43; for naphthalene 0.18. These values correspond satisfactorily with the data of other authors.

IthankM. N. Alentsev and M. D. Galanin for their valuable advice and help during the completion of this work.

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Translated by L. A. D'Asaro 308

The Surface Tension of Liquid He³ in the Region of Very LowTemperatures (1.0-0.35°K)

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J. Exper. Theoret. Phys. USSR 29, 899-900 (December, 1955)

 ${f E}$ ARLIER 1 we discussed our measurements of the surface tension of liquid He 3 from the criti-

cal temperature to 0.93 % by the method of capillary rise⁴. In the present work we have continued the measurements to lower temperatures using our earlier capillaries with diameters of 0.360, 0.224, and 0.188 mm. By pumping on He³ vapor, we were able to carry the measurements to 0.35°K.

The temperature of liquid He³ was determined by the vapor pressure of He³, measured with a Mcleod gauge. In the temperature interval 0.5-1.0°K we used the temperature scale of Sydoriak and Roberts (see reference 2). Below 0.5°K, the temperature was calculated from the vapor pressure, using an empirical formula given in reference 2. The pressure of the saturated vapor of He³ at a temperature of 0.35°K was 0.008 mm Hg. To measure the position of the level in the capillary we used a cathetometer, KM-5, which enabled us to read the height with an accuracy of 0.001 mm. However, as in the previous work, the actual accuracy of the reading was no better than several hundredths of a millimeter, since the observation of the meniscus was made through eight dewar vessel walls. The shape of the meniscus did not vary, and the guide line was adjusted for some average position of the division line. The error, introduced by the deformations of the heterogeneous glass, and noted by us earlier¹, was no greater than $\pm 2\%$.

The measurement proceeded in the following way. After the temperature had reached the lowest point, around 0.3° K, the pumping was stopped and the liquid slowly warmed up. In the warming process, regular measurements were taken of the vapor pressure of He³ and the height of rise of the liquid in the capillary. The warm-up time from 0.3° to 1.0° K took approximately 2.5 hours.

The calculation of the coefficient of surface tension was made according to the equation $\alpha = \rho g r h/2$, where the value of ρ , the density of liquid He³, was taken from the paper of E. C. Kerr³. The results obtained in the temperature interval 0.35-1.0°K, averaged for the three capillaries, are shown in Fig. 1 together with earlier data¹. Also shown are the results obtained by others^{4,5}. As we can see from the figure, the results from all three papers are in satisfactory agreement. (The old data have been recalculated using the new densities³.)

The error in determining the value of the surface tension in the region of very low temperature is approximately 5%. The maximum scatter of the individual points is shown by the arrows on the figure.

As we can see from the graph, the surface tension of μe^3 below 0.6°K approaches its limiting value equal to 0.152 dyne/cm. This value is 2.3 times smaller than the limiting value α_0 for He⁴, which is assumed to be equal to 0.35 dyne/cm on

^{*} The cap was constructed by E. E. Bukke.